

Constitutive Equations for Polymers

References:

- FA Morrison, *Understanding Rheology*, Oxford (2001)
- RG Larson, *Constitutive Equations for Polymer Melts and Solutions*, Butterworths (1988)
- RB Bird, RC Armstrong, O Hassager, *Dynamics of Polymeric Liquids*, Vol. 1+2, Wiley (1987)
- PJ Carreau, D DeKee, and RP Chhabra, *Rheology of Polymeric Systems*, Hanser (1997)

I. Elastic $\underline{\underline{\tau}}(t) = -G\underline{\underline{C}}^{-1}(t', t)$

II. Viscous

A. Newtonian $\underline{\underline{\tau}}(t) = -\mu\underline{\underline{\dot{\gamma}}}(t)$

B. Generalized Newtonian $\underline{\underline{\tau}}(t) = -\eta(\underline{\underline{\dot{\gamma}}})\underline{\underline{\dot{\gamma}}}(t)$ (see Carreau et al.)

1. Power Law GNF $\eta = m\dot{\gamma}^{n-1}$

2. Bingham Plastic GNF
$$\eta(\dot{\gamma}) = \begin{cases} \infty & |\underline{\underline{\tau}}| \leq \tau_0 \\ \mu_0 + \frac{\tau_0}{\dot{\gamma}} & |\underline{\underline{\tau}}| > \tau_0 \end{cases}$$

3. Carreau-Yasuda GNF $\eta = \eta_\infty + (\eta_0 - \eta_\infty) \left[1 + (\dot{\gamma}\lambda)^a \right]^{\frac{n-1}{a}}$

4. Ellis GNF
$$\eta = \frac{\eta_0}{1 + \left| \frac{\tau}{\tau_0} \right|^{\alpha-1}}$$

5. DeKee GNF $\eta = \eta_1 e^{-\lambda\dot{\gamma}} + \eta_2 e^{-0.1\lambda\dot{\gamma}} + \eta_\infty$

6. Casson GNF $\sqrt{\tau} = \sqrt{\tau_0} + \sqrt{\eta_0\dot{\gamma}}$

7. Herschel-Bulkley Model $\eta = \frac{\tau_0}{\dot{\gamma}} + m\dot{\gamma}^{n-1}$

8. DeKee-Turcotte Model $\eta = \frac{\tau_0}{\dot{\gamma}} + \eta_1 e^{-\lambda\dot{\gamma}}$

III. Linear Viscoelastic $\underline{\underline{\tau}}(t) = -\int_{-\infty}^t G(t-t') \underline{\underline{\dot{\gamma}}}(t') dt'$

A. Maxwell
$$G(t-t') = \left[\frac{\eta_0}{\lambda} e^{-\frac{(t-t')}{\lambda}} \right]$$

B. Generalized Maxwell $G(t-t') = \left[\sum_{k=1}^N \frac{\eta_k}{\lambda_k} e^{-\frac{(t-t')}{\lambda_k}} \right]$

C. Jeffreys $G(t-t') = \left[\frac{\eta_0}{\lambda_1} \left(1 - \frac{\lambda_2}{\lambda_1} \right) e^{-\frac{(t-t')}{\lambda_1}} + 2 \frac{\eta_0 \lambda_2}{\lambda_1} \delta(t-t') \right]$

D. Generalized Jeffreys $G(t-t') = \sum_{k=1}^N \left[\frac{\eta_k}{\lambda_k} \left(1 - \frac{\Lambda_k}{\lambda_k} \right) e^{-\frac{(t-t')}{\lambda_k}} + 2 \frac{\eta_k \Lambda_k}{\lambda_k} \delta(t-t') \right]$

IV. Nonlinear Viscoelastic

A. Continuum Modeling (organized by type of equation)

1. Differential Models

a. Quasilinear models

i. Upper-Convected Maxwell

$$\underline{\underline{\tau}} + \lambda \overset{\nabla}{\underline{\underline{\tau}}} = -\eta \dot{\underline{\underline{\gamma}}}$$

ii. Upper-Convected Jeffreys (Oldroyd B)

$$\underline{\underline{\tau}} + \lambda \overset{\nabla}{\underline{\underline{\tau}}} = -\eta \left(\dot{\underline{\underline{\gamma}}} + \lambda_2 \overset{\nabla}{\underline{\underline{\dot{\gamma}}}} \right)$$

iii. Lower-Convected Maxwell

$$\underline{\underline{\tau}} + \lambda \overset{\Delta}{\underline{\underline{\tau}}} = -\eta \dot{\underline{\underline{\gamma}}}$$

iv. Lower-Convected Jeffreys (Oldroyd A)

$$\underline{\underline{\tau}} + \lambda \overset{\Delta}{\underline{\underline{\tau}}} = -\eta \left(\dot{\underline{\underline{\gamma}}} + \lambda_2 \overset{\Delta}{\underline{\underline{\dot{\gamma}}}} \right)$$

v. Corotational Maxwell

$$\underline{\underline{\tau}} + \lambda \overset{\circ}{\underline{\underline{\tau}}} = -\eta \dot{\underline{\underline{\gamma}}}$$

vi. Corotational Jeffreys

$$\underline{\underline{\tau}} + \lambda \overset{\circ}{\underline{\underline{\tau}}} = -\eta \left(\dot{\underline{\underline{\gamma}}} + \lambda_2 \overset{\circ}{\underline{\underline{\dot{\gamma}}}} \right)$$

b. Models (quasi)linear in stress, at most quadratic in rate of deformation

i. Reiner-Rivlin

$$\underline{\underline{\tau}} = - \left(\phi_1 (II_{\dot{\underline{\underline{\gamma}}}}, III_{\dot{\underline{\underline{\gamma}}}}) \dot{\underline{\underline{\gamma}}} + \phi_2 (II_{\dot{\underline{\underline{\gamma}}}}, III_{\dot{\underline{\underline{\gamma}}}}) \dot{\underline{\underline{\gamma}}}^2 \right)$$

ii. Second-Order Fluid

$$\underline{\underline{\tau}} = - \left(\eta_0 \dot{\underline{\underline{\gamma}}} - \frac{\Psi_1^0}{2} \overset{\nabla}{\underline{\underline{\dot{\gamma}}}} + \Psi_2^0 \dot{\underline{\underline{\gamma}}} \cdot \dot{\underline{\underline{\gamma}}} \right)$$

iii. Oldroyd 8-constant (frame-invariance considerations) (see Bird, et al. Vol. 1)

- iv. Oldroyd 4-constant (frame-invariance considerations) (see Bird, et al. Vol. 1)
- v. Johnson-Segalman (nonaffine internal slip; see Larson)
- vi. Phan-Thien Tanner (nonaffine internal slip; see Larson)

c. Models quadratic in stress

- i. Giesekus model (anisotropic drag)

$$\underline{\underline{\tau}} + \lambda \overset{\nabla}{\underline{\underline{\tau}}} - \frac{\alpha \lambda}{\eta_0} \underline{\underline{\tau}} \cdot \underline{\underline{\tau}} = -\eta_0 \underline{\underline{\dot{\gamma}}}$$

- ii. Leonov model (leakage of elastic strain; see Larson)

2. Integral Models

a. Based on quasilinear models

- i. Lodge model

$$\underline{\underline{\tau}}(t) = - \int_{-\infty}^t \frac{\eta_0}{\lambda^2} e^{-\frac{(t-t')}{\lambda}} \underline{\underline{C}}^{-1}(t', t) dt'$$

- ii. Lodge rubberlike liquid

$$\underline{\underline{\tau}}(t) = - \int_{-\infty}^t \frac{\partial G(t-t')}{\partial t'} \underline{\underline{C}}^{-1}(t', t) dt'$$

- iii. Integral UC Jeffreys (Oldroyd B) model (same as Lodge Rubberlike Liquid with $G(t-t')$ given above under IIC)

- iv. Cauchy-Maxwell Model

$$\underline{\underline{\tau}}(t) = + \int_{-\infty}^t \frac{\eta_0}{\lambda^2} e^{-\frac{(t-t')}{\lambda}} \underline{\underline{C}}(t, t') dt'$$

- v. Integral LC Jeffreys (Oldroyd A) model (given below, with $G(t-t')$ given above under IIC)

$$\underline{\underline{\tau}}(t) = + \int_{-\infty}^t \frac{\partial G(t-t')}{\partial t'} \underline{\underline{C}}(t, t') dt'$$

b. General nonaffine motion models

- i. K-BKZ models (rubber elasticity theory; no molecular assumptions)

$$\underline{\underline{\tau}}(t) = + \int_{-\infty}^t M(t-t') \left(2 \frac{\partial U}{\partial I_2} \underline{\underline{C}} - 2 \frac{\partial U}{\partial I_1} \underline{\underline{C}}^{-1} \right) dt'$$

- ii. Rivlin-Sawyers models

$$\underline{\underline{\tau}}(t) = + \int_{-\infty}^t M(t-t') \left(\Phi_2(I_1, I_2) \underline{\underline{C}} - \Phi_1(I_1, I_2) \underline{\underline{C}}^{-1} \right) dt'$$

- B. Molecular Modeling (organized by type of model)
 - 1. Temporary Networks (see Larson)
 - a. Green-Tobolsky (UCM)
 - b. Yamamoto (breakage probability depends on chain extension)
 - c. Phan-Thien Tanner (Yamamoto with different closure)
 - 2. Bead Spring (see Bird, et al. Vol. 2)
 - a. Dumbbell (UCM)
 - b. Rouse (Generalized UCM)
 - c. Zimm
 - 3. Reptation (see Larson, current literature)
 - a. Doi-Edwards (RS or K-BKZ)
 - b. Refinements on DE (tube-length fluctuations, constraint release, chain stretching)
 - c. Pom-pom (branched polymers)