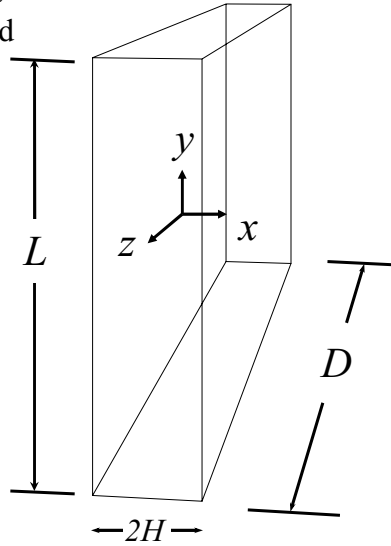
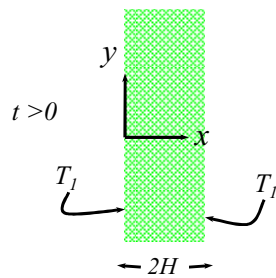


Example 1: Unsteady Heat Conduction in a Finite-sized solid

- The slab is tall and wide, but of thickness $2H$
- Initially at T_o
- at time $t = 0$ the temperature of the sides is changed to T_i



© Faith A. Morrison, Michigan Tech U.

Unsteady State Heat Transfer

Use same microscopic energy balance eqn as before.

$$\underbrace{\rho \hat{C}_p}_{\text{rate of change}} \left(\underbrace{\frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T}_{\text{convection}} \right) = \underbrace{k \nabla^2 T}_{\text{conduction (all directions)}} + \underbrace{S}_{\text{source (energy generated per unit volume per time)}}$$

see handout for component notation

© Faith A. Morrison, Michigan Tech U.

Microscopic Energy Equation in Cartesian Coordinates

$$\frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} = \frac{k}{\rho \hat{C}_p} \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + \frac{S}{\rho \hat{C}_p}$$

$$\alpha \equiv \frac{k}{\rho \hat{C}_p} = \text{thermal diffusivity}$$

what are the boundary conditions? initial conditions?

© Faith A. Morrison, Michigan Tech U.

Unsteady State Heat Conduction in a Finite Slab

$$\frac{\partial T}{\partial t} = \frac{k}{\rho \hat{C}_p} \left(\frac{\partial^2 T}{\partial x^2} \right) = \alpha \left(\frac{\partial^2 T}{\partial x^2} \right)$$

Initial condition: $t = 0, T = T_o \forall x$

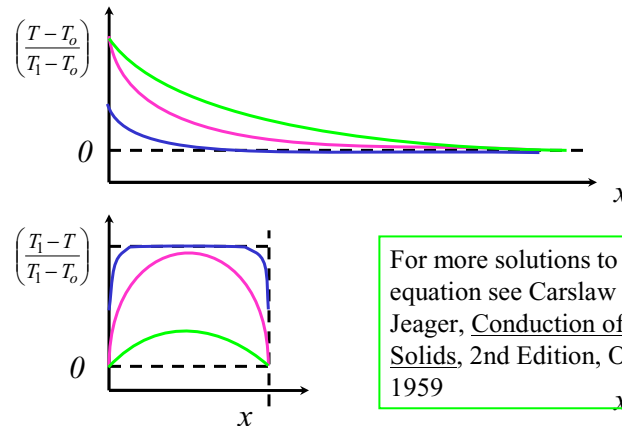
Boundary conditions:

$$\left. \begin{array}{l} x = 0, \quad T = T_1 \\ x = 2H, \quad T = T_1 \end{array} \right\} \forall t > 0$$

© Faith A. Morrison, Michigan Tech U.

Q: How can two completely different situations give the same governing equation?

A: The boundary conditions make all the difference



For more solutions to this equation see Carslaw and Jeager, Conduction of Heat in Solids, 2nd Edition, Oxford, 1959

© Faith A. Morrison, Michigan Tech U.

Unsteady State Heat Conduction in a Finite Slab:
solution by separation of variables

Let $Y \equiv \left(\frac{T_1 - T}{T_1 - T_o} \right)$ $\frac{\partial Y}{\partial t} = \alpha \left(\frac{\partial^2 Y}{\partial x^2} \right)$

Guess: $Y = X(x)\Theta(t)$

Initial condition:

$t = 0, T = T_o \forall x \Rightarrow Y = 1$

Boundary conditions:

$\left. \begin{matrix} x = 0, & T = T_1 \Rightarrow Y = 0 \\ x = 2H, & T = T_1 \Rightarrow Y = 0 \end{matrix} \right\} \forall t > 0$

© Faith A. Morrison, Michigan Tech U.

Unsteady State Heat Conduction in a Finite Slab: soln by separation of variables

$$Y = X(x)\Theta(t) \quad \frac{\partial Y}{\partial t} = \alpha \left(\frac{\partial^2 Y}{\partial x^2} \right)$$

$$\frac{\partial Y}{\partial t} = \frac{\partial}{\partial t} (X(x)\Theta(t)) = X(x) \frac{d\Theta(t)}{dt}$$

$$\frac{\partial Y}{\partial x} = \frac{\partial}{\partial x} (X(x)\Theta(t)) = \frac{dX(x)}{dx} \Theta(t)$$

$$\frac{\partial^2 Y}{\partial x^2} = \frac{d^2 X(x)}{dx^2} \Theta(t)$$

© Faith A. Morrison, Michigan Tech U.

Unsteady State Heat Conduction in a Finite Slab: soln by separation of variables

$$\frac{\partial Y}{\partial t} = \alpha \left(\frac{\partial^2 Y}{\partial x^2} \right)$$

Substituting: $X(x) \frac{d\Theta(t)}{dt} = \alpha \frac{d^2 X(x)}{dx^2} \Theta(t)$

$$\underbrace{\frac{1}{\Theta(t)} \frac{d\Theta(t)}{dt}}_{\text{function of time only}} = \alpha \underbrace{\frac{1}{X(x)} \frac{d^2 X(x)}{dx^2}}_{\text{function of position (x) only}} \Rightarrow = \lambda \text{ constant}$$

© Faith A. Morrison, Michigan Tech U.

Unsteady State Heat Conduction in a Finite Slab: soln by separation of variables

Separates into two ordinary differential equations:

$$\frac{1}{\Theta(t)} \frac{d\Theta(t)}{dt} = \lambda$$
$$\alpha \frac{1}{X(x)} \frac{d^2 X(x)}{dx^2} = \lambda$$

Solve.

Apply BCs.

Apply ICs.

© Faith A. Morrison, Michigan Tech U.

Temperature Profile for Unsteady State Heat Conduction in a Finite Slab

$$\left(\frac{T_1 - T}{T_1 - T_0} \right) = \frac{4}{\pi} \left\{ e^{-\frac{\pi^2 \alpha t}{4H^2}} \sin \frac{\pi x}{2H} + \frac{1}{3} e^{-\frac{3^2 \pi^2 \alpha t}{4H^2}} \sin \frac{3\pi x}{2H} + \frac{1}{5} e^{-\frac{5^2 \pi^2 \alpha t}{4H^2}} \sin \frac{5\pi x}{2H} + \dots \right\}$$

Geankoplis 4th ed., eqn 5.3-6, p363

© Faith A. Morrison, Michigan Tech U.

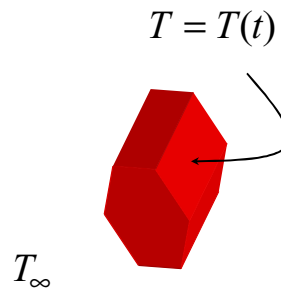
Unsteady Macroscopic Energy Balances

Example 3: If a piece of steel at $T = T_o$ is dropped into a large reservoir of fluid at T_∞ , what is the temperature of the steel as a function of time?

$k = \text{large}$, which means that there is no resistance to heat transfer in the steel

Therefore, we are NOT calculating a temperature profile

Use Macroscopic Energy Balance

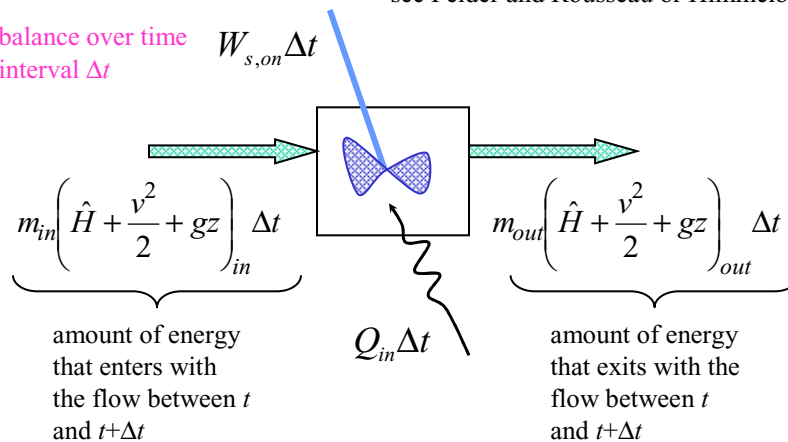


© Faith A. Morrison, Michigan Tech U.

Unsteady Macroscopic Energy Balances

see Felder and Rousseau or Himmelblau

balance over time interval Δt



© Faith A. Morrison, Michigan Tech U.

Unsteady Macroscopic Energy Balances

accumulation = input - output

$$\frac{d(U_{sys} + E_{k,sys} + E_{p,sys})}{dt} = -\Delta H - \Delta E_p - \Delta E_k + Q_{in} + W_{s,on}$$

Our case, no shafts

For negligible changes in E_p and E_k ,
and no phase change,
and single-input, single-output system,

$$\frac{dU_{sys}}{dt} = M_{sys} C_v \frac{dT_{sys}}{dt}$$

$$\Delta H = m C_p (T_{out} - T_{in})$$

Our case, no flow

© Faith A. Morrison, Michigan Tech U.

Unsteady Macroscopic Energy Balances

$$M_{sys} C_v \frac{dT_{sys}}{dt} = m C_p (T_{in} - T_{sys}) + Q_{in} + W_{s,on}$$

Our case, no flow

Our case, no shafts

- For negligible changes in E_p and E_k
- and no phase change
- and single-input, single-output system
- assume $T_{out} = T_{sys}$

© Faith A. Morrison, Michigan Tech U.

Unsteady Macroscopic Energy Balance
Applied to cooling steel part:

$$M_{\text{sys}} C_v \frac{dT_{\text{sys}}}{dt} = Q_{\text{in}}$$

The temperature changes in the slug are due to the heat loss

The heat loss depends on the heat-transfer coefficient from the slug to the environment

$$Q_{\text{in}} = hA(T - T_{\infty})$$