

CM3110

Transport Processes and Unit Operations I

Fluid Mechanics
*Non-Newtonian fluids –
 An Introduction*

The Weissenberg effect is when a viscoelastic, non-Newtonian fluid will climb a rotating shaft.

<https://www.youtube.com/watch?v=npZzlgKjs0I>



Photo by Carlos Arango Sabogal, U. Wisconsin, Madison

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Intro to Non-Newtonian Fluid Mechanics

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Continuum Modeling—Newtonian

fluid particle pathlines

control volume

- μ (viscosity) constant
- ρ (density) constant often a good assumption
- $\underline{\tau} = \mu(\nabla v + (\nabla v)^T)$

The form of the function
 $\underline{\tau}(v)$ is known

—Viscosity differs for different materials
 —Is a function of T, perhaps P

This is MAJOR.
 Predictions seem to be right in a wide variety of situations (water, oil, air)

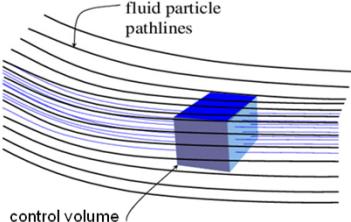
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Intro to Non-Newtonian Fluid Mechanics



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Continuum Modeling—Non-Newtonian



- μ (viscosity) NOT constant
- ρ (density) constant is often a good assumption
- $\underline{\tau} = \dots ?$ } The form of the function
 $\underline{\tau}(v)$ is NOT known

For non-Newtonian Fluids, we need a new form of $\underline{\tau}(v)$ that matches material observations.

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Intro to Non-Newtonian Fluid Mechanics



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When did we separate out non-Newtonian fluids?

Lecture 2:

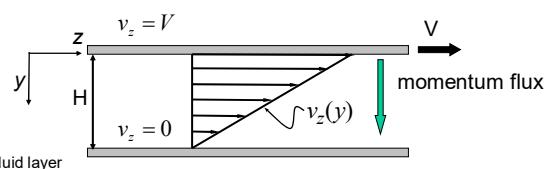
How do Fluids Behave?

Momentum Flux

Momentum (p) = mass * velocity

$$p = m\underline{v} \quad \text{vectors}$$

top plate has momentum, and it transfers this momentum to the top layer of fluid



Each fluid layer transfers the momentum downward

Viscosity determines the magnitude of momentum flux

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Intro to Non-Newtonian Fluid Mechanics

The graph shows a linear relationship between shear stress $\tilde{\tau}_{21}$ on the vertical axis and shear rate $\dot{\gamma} = \left| \frac{\partial v_1}{\partial x_2} \right|$ on the horizontal axis. A straight blue line passes through the origin, labeled "Newtonian". A right-angled triangle is drawn on the line to indicate its slope, which is labeled "slope = μ , viscosity".

Newtonian Fluids

Newton's Law of Viscosity
(unidirectional flow)

$$\tilde{\tau}_{21} = \mu \left(\frac{\partial v_1}{\partial x_2} \right)$$

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Non-Newtonian Fluids

Non-Newtonian behavior
(unidirectional flow)

The graph shows three curves starting from a common origin. A black line represents "Bingham plastic" behavior, showing an initial yield stress $\tilde{\tau}_0$ before it begins to flow. A blue line represents "Newtonian" behavior, starting at the origin. A red curve represents "shear-thickening or dilatant" behavior, where the shear stress increases with shear rate. Labels indicate "shear-thinning or pseudoplastic" for the red curve and "shear-thickening or dilatant" for the red curve.

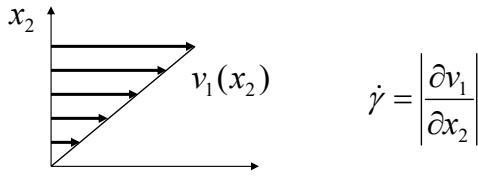
$\tilde{\tau}_{21} = \text{function of } \left| \frac{\partial v_1}{\partial x_2} \right|$

Many different behaviors are observed.

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How can we define viscosity for Non-Newtonian Fluids?

- Perform this:



- Measure $\tilde{\tau}_{21}$

- Calculate this:

$$\eta \equiv \frac{\tilde{\tau}_{21}}{\dot{\gamma}}$$

Non-Newtonian viscosity
(unidirectional flow)

(NOTE on coordinate system: Viscosity definition is written for shear flow in x_1 direction and gradient in x_2 direction)

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Typical polymeric behavior

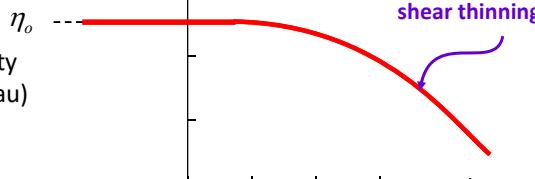
Non-Newtonian behavior
(unidirectional flow)

The changes in viscosity with shear rate are so large they must be plotted log-log

$\log \eta$

Viscosity (Greek letter eta)

zero-shear viscosity
(Newtonian plateau)



$$\dot{\gamma} \equiv \left| \frac{\partial v_1}{\partial x_2} \right|$$

shear rate ("gamma dot")

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In addition, for many polymers there are shear-induced **NORMAL** (perpendicular) forces.

force on 2-surface in 2-direction

$$\iint_S (\tilde{\tau}_{22} - p) dS$$

force on 2-surface in 1-direction

$$\iint_S \tilde{\tau}_{21} dS$$



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How to deal with this?

Recall, for Newtonian fluids:

$\underline{\underline{\tilde{\tau}}} = \mu \dot{\underline{\underline{v}}} = \mu (\nabla \underline{v} + (\nabla \underline{v})^T)$	Newtonian Constitutive Equation
$\begin{pmatrix} \tilde{\tau}_{xx} & \tilde{\tau}_{xy} & \tilde{\tau}_{xz} \\ \tilde{\tau}_{yx} & \tilde{\tau}_{yy} & \tilde{\tau}_{yz} \\ \tilde{\tau}_{zx} & \tilde{\tau}_{zy} & \tilde{\tau}_{zz} \end{pmatrix}_{xyz} = \mu \begin{pmatrix} 2 \frac{\partial v_x}{\partial x} & \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} & \frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} \\ \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} & 2 \frac{\partial v_y}{\partial y} & \frac{\partial v_z}{\partial y} + \frac{\partial v_y}{\partial z} \\ \frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} & \frac{\partial v_z}{\partial y} + \frac{\partial v_y}{\partial z} & 2 \frac{\partial v_z}{\partial z} \end{pmatrix}_{xyz}$	
$\underline{v} = \begin{pmatrix} 0 \\ v_z \end{pmatrix}_{xyz} \Rightarrow \tilde{\tau}_{xz} = \mu \left(\frac{dv_z}{dx} \right)$	<i>Newton's law of viscosity is a special case of the Newtonian Constitutive equation. (Unidirectional flow)</i>

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$$\tilde{\tau} = \mu \dot{\gamma} = \mu (\nabla \underline{v} + (\nabla \underline{v})^T)$$

**Non-Newtonian
Newtonian
Constitutive
Equation**

$$\begin{pmatrix} \tilde{\tau}_{xx} & \tilde{\tau}_{xy} & \tilde{\tau}_{xz} \\ \tilde{\tau}_{yx} & \tilde{\tau}_{yy} & \tilde{\tau}_{yz} \\ \tilde{\tau}_{zx} & \tilde{\tau}_{zy} & \tilde{\tau}_{zz} \end{pmatrix}_{xyz} = \mu \begin{pmatrix} 2 \frac{\partial v_x}{\partial x} & \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} & \frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} \\ \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} & 2 \frac{\partial v_y}{\partial y} & \frac{\partial v_z}{\partial y} + \frac{\partial v_y}{\partial z} \\ \frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} & \frac{\partial v_z}{\partial y} + \frac{\partial v_y}{\partial z} & 2 \frac{\partial v_z}{\partial z} \end{pmatrix}_{xyz}$$

**Generalized
Newtonian Fluid
(non-Newtonian)**

$\eta(\dot{\gamma})$

We pick the form of this function that works best with our data.

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Power-Law Model

Non-Newtonian behavior
(unidirectional flow)

$$\tilde{\tau}_{21} = \text{function of } \left(\frac{\partial v_1}{\partial x_2} \right)$$

$\eta(\dot{\gamma}) = m \dot{\gamma}^{n-1}$

(does **not** model normal stresses)

$$\tilde{\tau}_{21} = m \left| \frac{dv_1}{dx_2} \right|^{n-1} \frac{dv_1}{dx_2}$$

m or K = consistency index ($m = \mu$ for Newtonian)
 n = power-law index ($n = 1$ for Newtonian)

$\dot{\gamma} \equiv \left| \frac{\partial v_1}{\partial x_2} \right| = \text{shear rate}$

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What does the power-law model predict for viscosity?

$$\tilde{\tau}_{21} = m \left| \frac{dv_1}{dx_2} \right|^{n-1} \frac{dv_1}{dx_2}$$

$$\eta \equiv \frac{\tilde{\tau}_{21}}{\dot{\gamma}} = \frac{\tilde{\tau}_{21}}{\left| \frac{dv_1}{dx_2} \right|} = m \left| \frac{dv_1}{dx_2} \right|^{n-1}$$

On a log-log plot, this would give
a straight line:

$$\log \eta = \log m + (n-1) \log \left| \frac{dv_1}{dx_2} \right|$$

$\underbrace{Y}_{Y} = \underbrace{B}_{B} + \underbrace{M}_{M} \underbrace{X}_{X}$

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Power-Law Fluid

Non-Newtonian
viscosity

$$\eta \equiv \frac{\tilde{\tau}_{21}}{\dot{\gamma}} \quad \dot{\gamma} = \left| \frac{\partial v_1}{\partial x_2} \right|$$

$$\eta = m \dot{\gamma}^{n-1}, \quad n > 1$$

shear thickening

$$\text{Newtonian } \eta = m \dot{\gamma}^{n-1}, \quad n = 1$$

$$\eta = m \dot{\gamma}^{n-1}, \quad n < 1$$

shear thinning

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Where do we use the power-law expression?

The Equation of Continuity and the Equation of Motion in Cartesian, cylindrical, and spherical coordinates

Continuity Equation, Cartesian coordinates

$$\frac{\partial \rho}{\partial t} + \left(v_x \frac{\partial \rho}{\partial x} + v_y \frac{\partial \rho}{\partial y} + v_z \frac{\partial \rho}{\partial z} \right) + \rho \left(\frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \right) = 0$$

Continuity Equation, cylindrical coordinates

$$\frac{\partial \rho}{\partial t} + \frac{1}{r} \frac{\partial(\rho r v_r)}{\partial r} + \frac{1}{r} \frac{\partial(\rho v_\theta)}{\partial \theta} + \frac{\partial(\rho v_z)}{\partial z} = 0 \quad \frac{\partial \rho}{\partial t} + \underline{v} \cdot \nabla \rho = -\rho (\nabla \cdot \underline{v})$$

Continuity Equation, spherical coordinates

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial(\rho r^2 v_r)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial(\rho v_\theta \sin \theta)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial(\rho v_\phi)}{\partial \phi} = 0$$

Equation of Motion for an incompressible fluid, 3 components in Cartesian coordinates

$$\begin{aligned} \rho \left(\frac{\partial v_x}{\partial t} + v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} + v_z \frac{\partial v_x}{\partial z} \right) &= -\frac{\partial P}{\partial x} + \left(\frac{\partial \tilde{\tau}_{xx}}{\partial x} + \frac{\partial \tilde{\tau}_{xy}}{\partial y} + \frac{\partial \tilde{\tau}_{xz}}{\partial z} \right) + \rho g_x \\ \rho \left(\frac{\partial v_y}{\partial t} + v_x \frac{\partial v_y}{\partial x} + v_y \frac{\partial v_y}{\partial y} + v_z \frac{\partial v_y}{\partial z} \right) &= -\frac{\partial P}{\partial y} + \left(\frac{\partial \tilde{\tau}_{xy}}{\partial x} + \frac{\partial \tilde{\tau}_{yy}}{\partial y} + \frac{\partial \tilde{\tau}_{yz}}{\partial z} \right) + \rho g_y \\ \rho \left(\frac{\partial v_z}{\partial t} + v_x \frac{\partial v_z}{\partial x} + v_y \frac{\partial v_z}{\partial y} + v_z \frac{\partial v_z}{\partial z} \right) &= -\frac{\partial P}{\partial z} + \left(\frac{\partial \tilde{\tau}_{xz}}{\partial x} + \frac{\partial \tilde{\tau}_{yz}}{\partial y} + \frac{\partial \tilde{\tau}_{zz}}{\partial z} \right) + \rho g_z \end{aligned}$$

The one
with "̃" is
for non-
Newtonian
fluids

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www.chem.mtu.edu/~fmorriso/cm310/Navier.pdf

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We have a handout for that:

The Power-Law, Generalized Newtonian Constitutive Equation in Cartesian, Cylindrical, and Spherical coordinates

Prof. Faith A. Morrison, Michigan Technological University

Cartesian Coordinates

$$\begin{aligned} \begin{pmatrix} \tilde{\tau}_{xx} & \tilde{\tau}_{xy} & \tilde{\tau}_{xz} \\ \tilde{\tau}_{yx} & \tilde{\tau}_{yy} & \tilde{\tau}_{yz} \\ \tilde{\tau}_{zx} & \tilde{\tau}_{zy} & \tilde{\tau}_{zz} \end{pmatrix}_{xyz} &= \eta \dot{\underline{v}} \\ \eta = m \left(\frac{1}{2} \cdot \text{sum of squares of each term in } \dot{\underline{v}} \right)^{\frac{n-1}{2}} &= m \left(\frac{1}{2} \sum_{p=1}^3 \sum_{j=1}^3 \dot{v}_p^2 \right)^{\frac{n-1}{2}} \\ \dot{\underline{v}} \equiv &\begin{pmatrix} 2 \frac{\partial v_x}{\partial x} & \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} & \frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} \\ \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} & 2 \frac{\partial v_y}{\partial y} & \frac{\partial v_y}{\partial z} + \frac{\partial v_z}{\partial y} \\ \frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} & \frac{\partial v_y}{\partial z} + \frac{\partial v_z}{\partial y} & 2 \frac{\partial v_z}{\partial z} \end{pmatrix}_{xyz} \end{aligned}$$

Cylindrical Coordinates

$$\begin{aligned} \begin{pmatrix} \tilde{\tau}_{rr} & \tilde{\tau}_{r\theta} & \tilde{\tau}_{rz} \\ \tilde{\tau}_{\theta r} & \tilde{\tau}_{\theta\theta} & \tilde{\tau}_{\theta z} \\ \tilde{\tau}_{zr} & \tilde{\tau}_{\theta z} & \tilde{\tau}_{zz} \end{pmatrix}_{r\theta z} &= \eta \dot{\underline{v}} \\ \eta = m \left(\frac{1}{2} \cdot \text{sum of squares of each term in } \dot{\underline{v}} \right)^{\frac{n-1}{2}} &= m \left(\frac{1}{2} \sum_{p=1}^3 \sum_{j=1}^3 \dot{v}_p^2 \right)^{\frac{n-1}{2}} \\ \dot{\underline{v}} \equiv &\begin{pmatrix} 2 \frac{\partial v_r}{\partial r} & r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} & \frac{\partial v_r}{\partial z} + \frac{\partial v_z}{\partial r} \\ r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} & 2 \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r}{r} & \frac{1}{r} \frac{\partial v_z}{\partial \theta} + \frac{\partial v_\theta}{\partial z} \\ \frac{\partial v_r}{\partial z} + \frac{\partial v_z}{\partial r} & \frac{1}{r} \frac{\partial v_z}{\partial \theta} + \frac{\partial v_\theta}{\partial z} & 2 \frac{\partial v_z}{\partial z} \end{pmatrix}_{r\theta z} \end{aligned}$$

<http://www.chem.mtu.edu/~fmorriso/cm310/stpl.pdf>

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Where do we use the power-law expression?

e.g., Poiseuille flow in a tube:

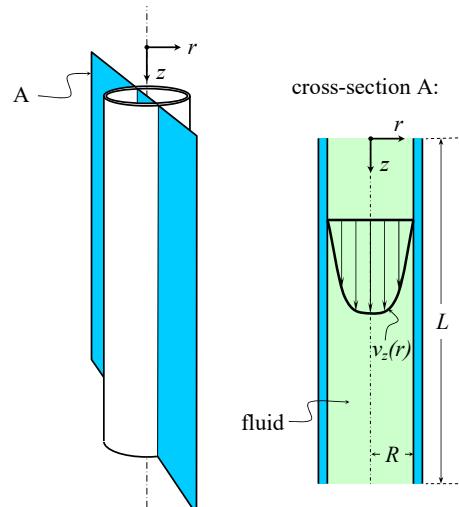
$$\tilde{\tau}_{rz} = \mu \left(\frac{dv_z}{dr} \right) \quad \text{Newtonian}$$

$$\tilde{\tau}_{rz} = m \left| \frac{dv_z}{dr} \right|^{n-1} \frac{dv_z}{dr} \quad \text{non-Newtonian, power-law}$$

$$\tilde{\tau}_{rz} = \left(\frac{L\rho g + (P_o - P_L)}{2L} \right) r \quad \Rightarrow \quad \text{solve for } v_z(r)$$

1-direction = r
2-direction = z

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EXAMPLE III: Pressure-driven flow of a Power-law fluid in a tube

- steady state
- incompressible
- well developed
- long tube

Calculate
velocity and
stress profiles

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Non-Newtonian behavior
(unidirectional flow)

Calculate the velocity field for Pressure-driven flow of a power-law (PL) fluid:

$$\tilde{\tau}_{rz} = m \left| \frac{dv_z}{dr} \right|^{n-1} \frac{dv_z}{dr} = \left(\frac{L\rho g + (P_o - P_L)}{2L} \right) r$$

$\underbrace{\phantom{m \left| \frac{dv_z}{dr} \right|^{n-1}}}_{\partial v_z / \partial r} \quad \underbrace{\phantom{\left(\frac{L\rho g + (P_o - P_L)}{2L} \right) r}}_{\equiv \alpha}$

$$= -\frac{\partial v_z}{\partial r} \quad \forall r \quad \equiv \alpha$$

$$m \left(-\frac{\partial v_z}{\partial r} \right)^{n-1} \frac{\partial v_z}{\partial r} = -m \left(-\frac{\partial v_z}{\partial r} \right)^n = \alpha r$$

Solve for $v_z(r)$

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Boundary Conditions: ?

(same as before in the Newtonian case)

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Non-Newtonian behavior
(unidirectional flow)

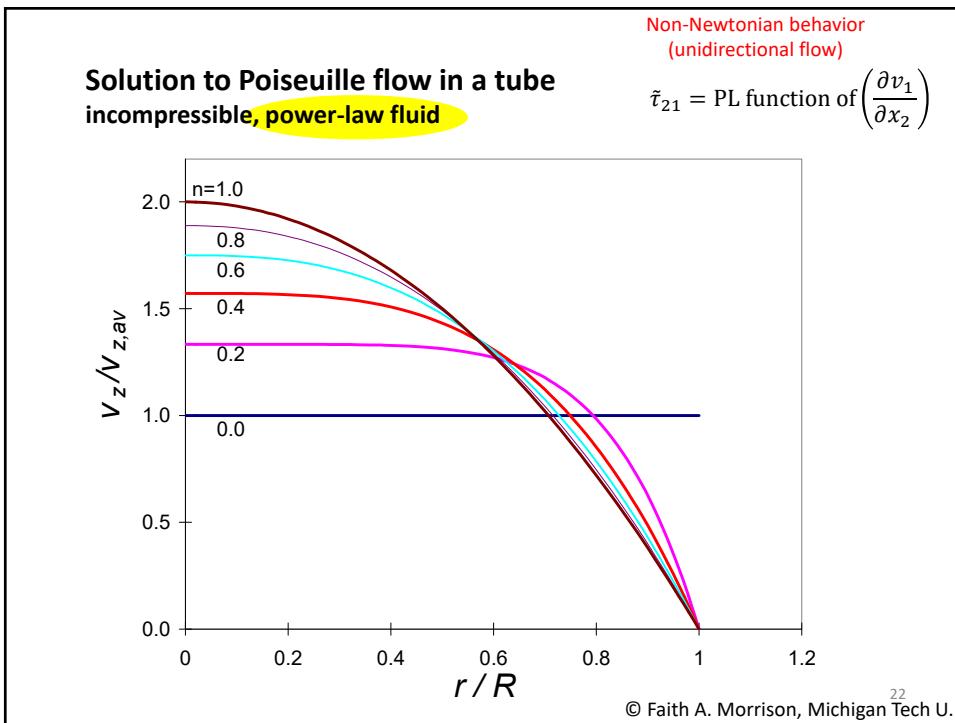
Velocity field

Poiseuille flow of a power-law fluid:

$$v_z(r) = \left(\frac{R(L\rho g + P_o - P_L)}{2Lm} \right)^{\frac{1}{n}} \left(\frac{R}{\frac{1}{n} + 1} \right) \left(1 - \left(\frac{r}{R} \right)^{\frac{1}{n}+1} \right)$$

$$\langle v \rangle = \frac{1}{\pi R^2} \int_0^{2\pi} \int_0^R v_z(r) r dr d\theta = R \left(\frac{n}{1+3n} \right) \left[\frac{R(P_o - P_L)}{2mL} \right]^{\frac{1}{n}}$$

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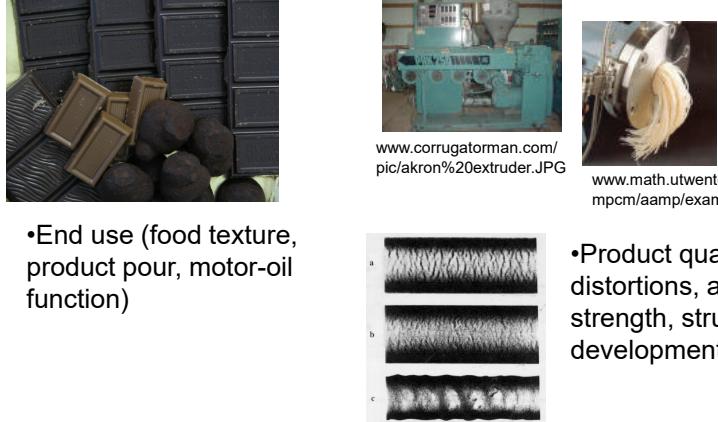


Non-Newtonian behavior (all flows)
 $\tilde{\tau}_{21}$ = nonlinear function of $\nabla \mathcal{V}$

Rheology (Non-Newtonian Fluid Mechanics)

Rheology affects:

- Processing (design, costs, production rates)
- End use (food texture, product pour, motor-oil function)
- Product quality (surface distortions, anisotropy, strength, structure development)



www.corrugatorman.com/pic/akron%20extruder.JPG

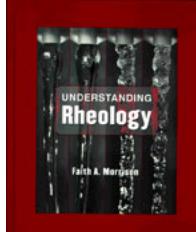
www.math.utwente.nl/mpcm/aamp/examples.html

Pomar et al. JNNFM
54 143 1994

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Rheology (Non-Newtonian Fluid Mechanics)

At Michigan Tech:
CM4650 Polymer Rheology (Even years spring)
www.chem.mtu.edu/~fmorriso/cm4650/cm4650.html



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Next:

CM3110
Transport I
Part I: Fluid Mechanics

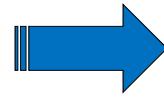
More Complicated Flows
(Dimensional Analysis,
rough pipes, hydraulic
diameter, porous media)

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